



Study of Intermetallic Nanostructures for Light-Water Reactor Materials

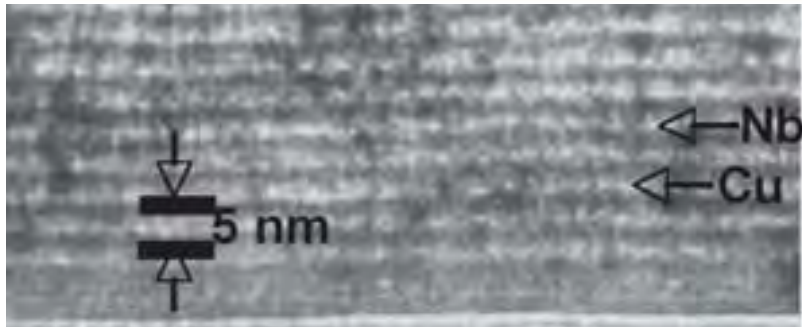
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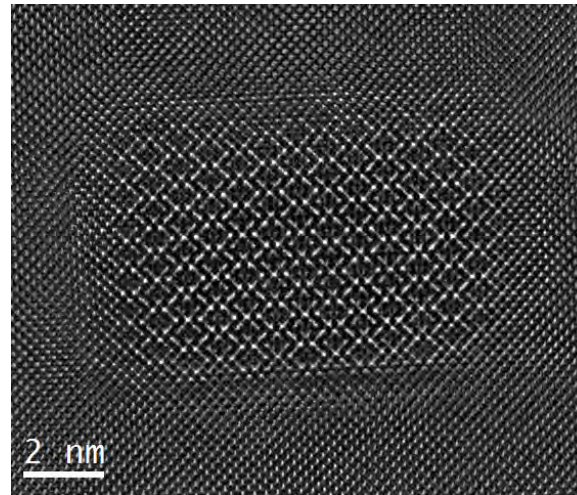
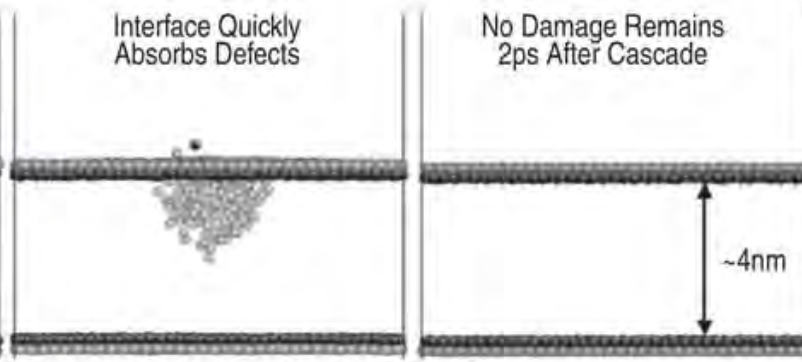
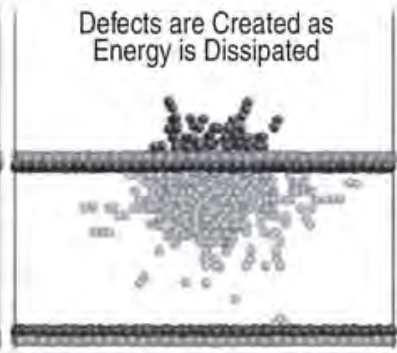
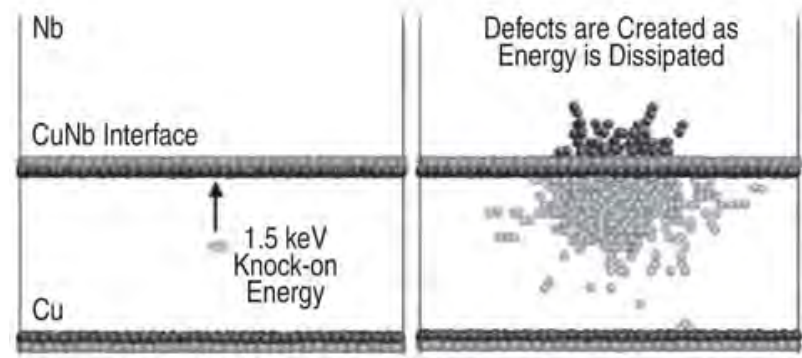
OUTLINE

- **The idea/concept** (pursuing interfaces as sinks for defects)
- **Requirements for a new intermetallic forming steel in nuclear environment** (simple manufacturing, reduced cost, good properties)
- **Aging and heat treatment evaluation of the intermetallic forming steels**
- **Radiation damage study. The beginning.**
- **Theoretical work** (broad development of tools with specific directed applications: algorithms, potentials, simulations)
- **Future plans**

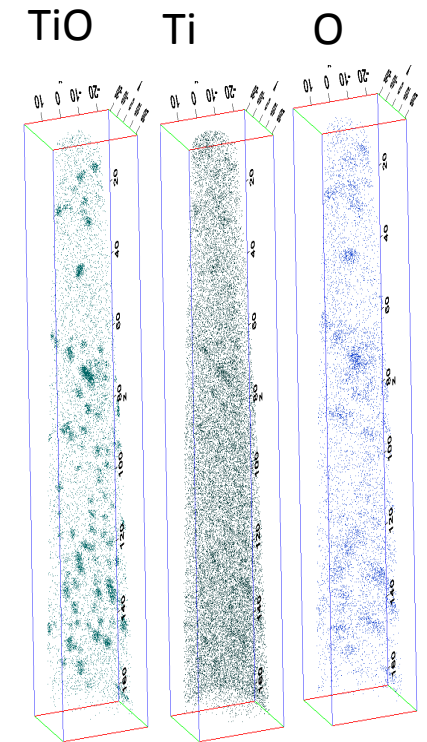
The idea



Radiation tolerant materials are materials with high interface density.
Cu/Nb layered interfaces or ODS alloys have been studied for decades.



HR TEM image of MA957

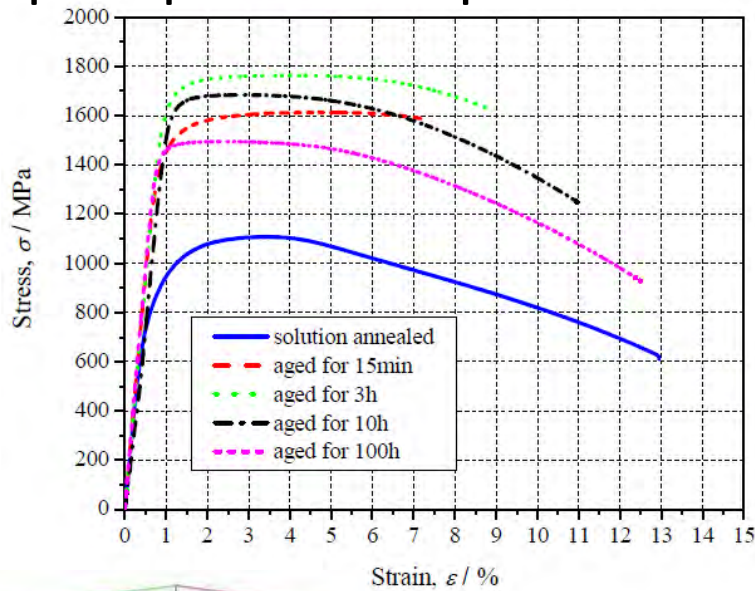


AtomProbeTomography data on MA957

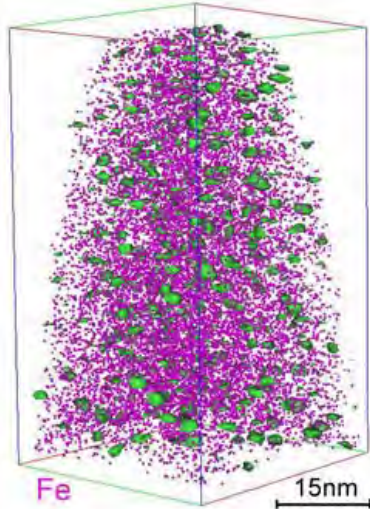
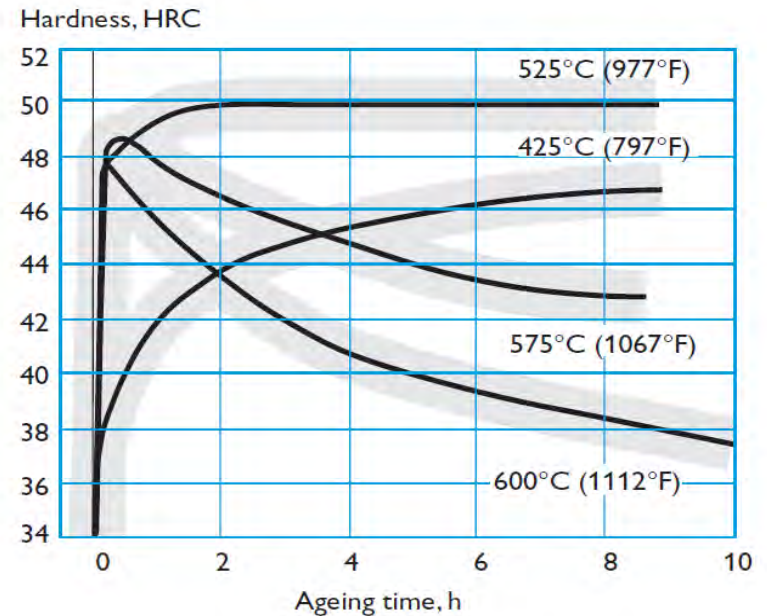
What about intermetallic's??

Intermetallic forming steels

Mar(tensite)aging steels are known to form NiAl and Ni₃Al precipitates at specific heat treatments leading to extraordinary properties



properties



Green: NiAl

Red: Fe

Also Cr

APT measurement by
M. Schober et al.
Univ. of Leoben, Austria
NiAl precipitates shown
green

Corrax

| | | | | | | | |
|--------------------|-----------------------------|-----|-----|------|-----|-----|-----|
| Typical analysis % | C | Si | Mn | Cr | Ni | Mo | Al |
| | 0,03 | 0,3 | 0,3 | 12,0 | 9,2 | 1,4 | 1,6 |
| Delivery condition | Solution treated to ~34 HRC | | | | | | |
| Colour code | Black/grey | | | | | | |

Requirements for new materials

Stability under radiation

→ radiation damage studies

Stability at high elevated temperatures for an extended period of time.

→ Aging studies

Understanding the precipitation in relevant environments

→ Aging studies

Reasonable strength,

→ Mechanical properties

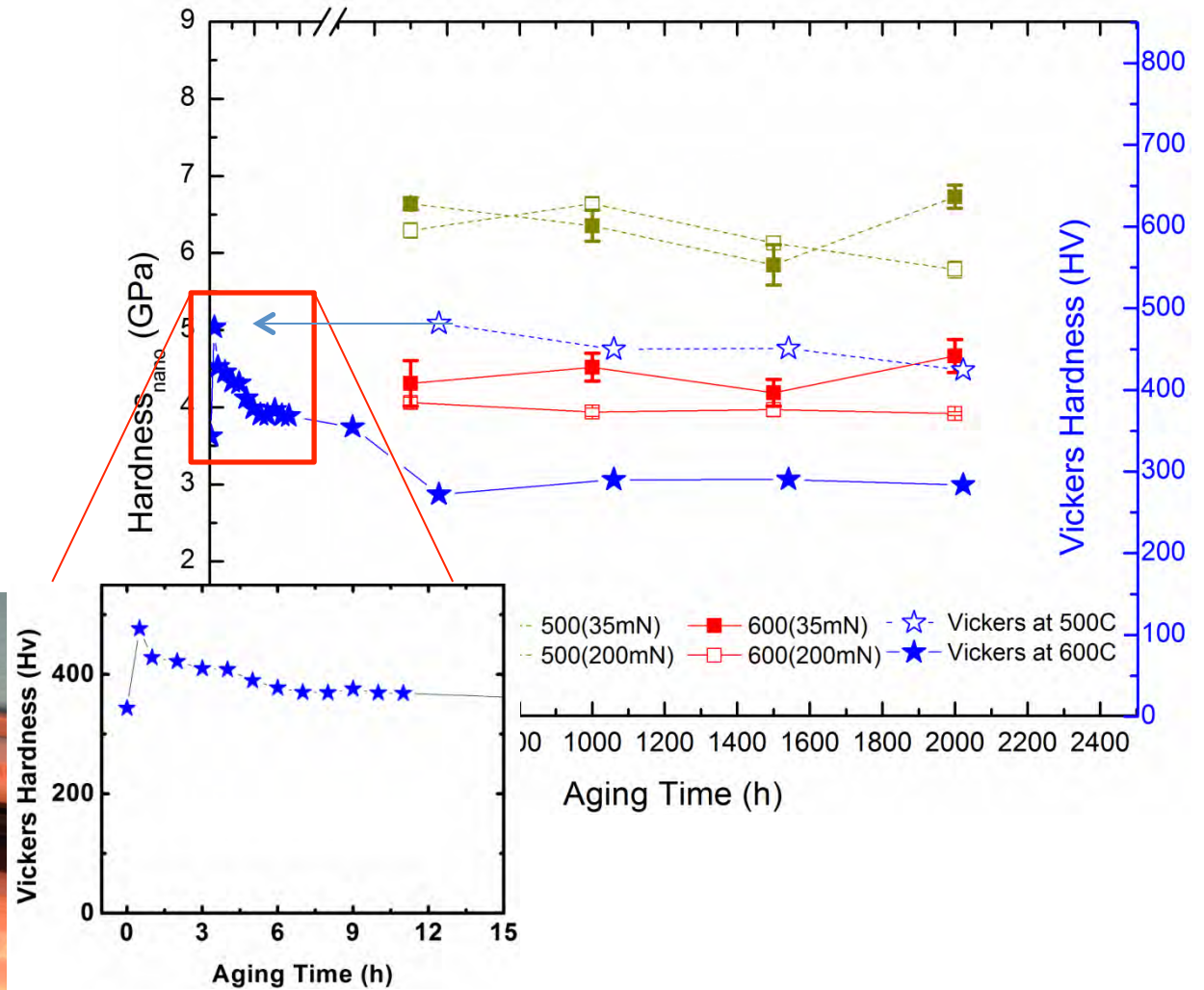
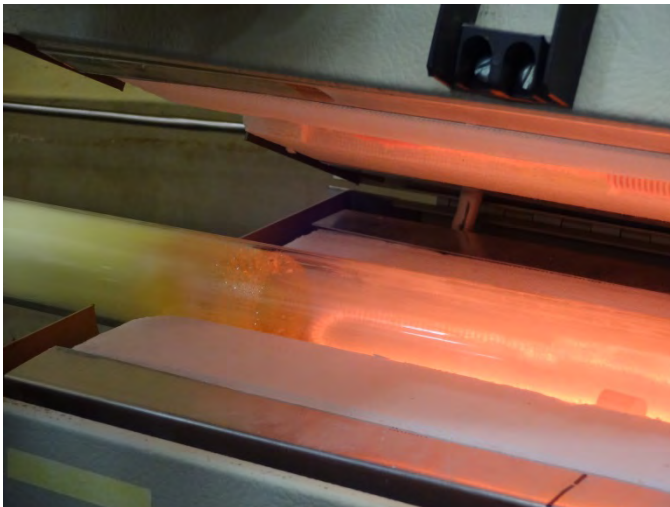
Reasonable elongation/ ductility

→ Mechanical properties

Reasonably low activation in a neutron environment

→ Co free materials

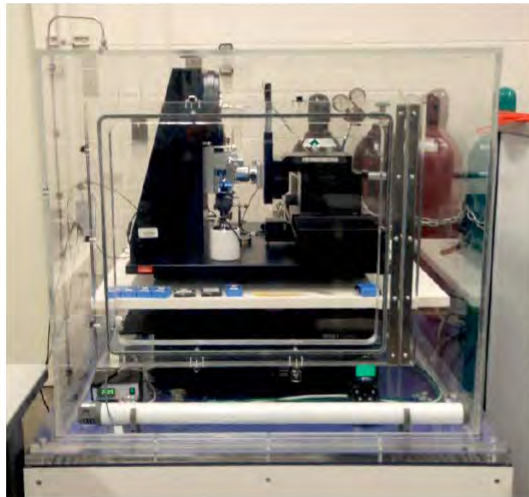
Experimental work performed on long term aging



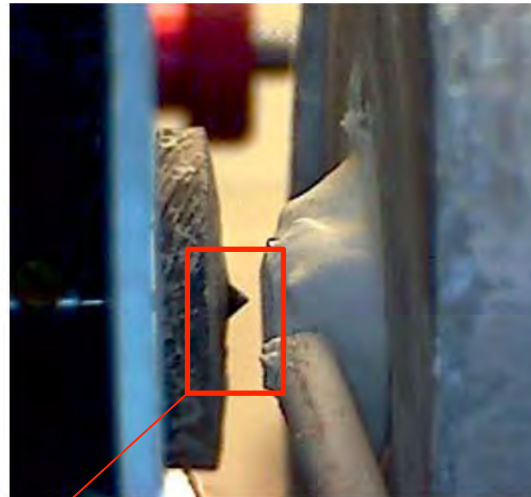
35mN and 200mN nanoindentation results
Vickers measured with 200g load (HV0.2)

Experimental work performed on heat treatments

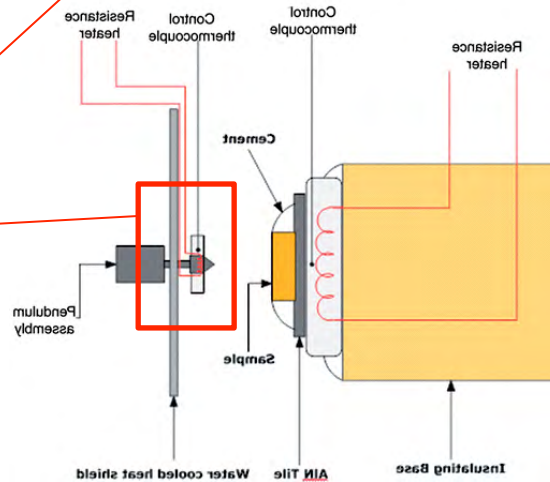
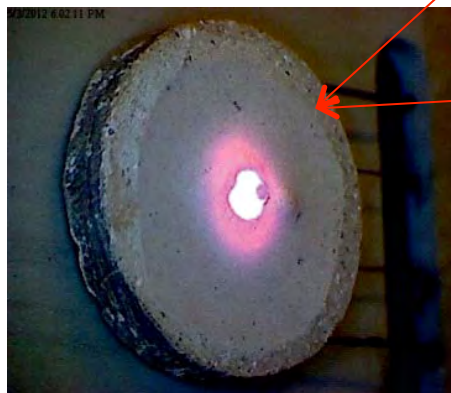
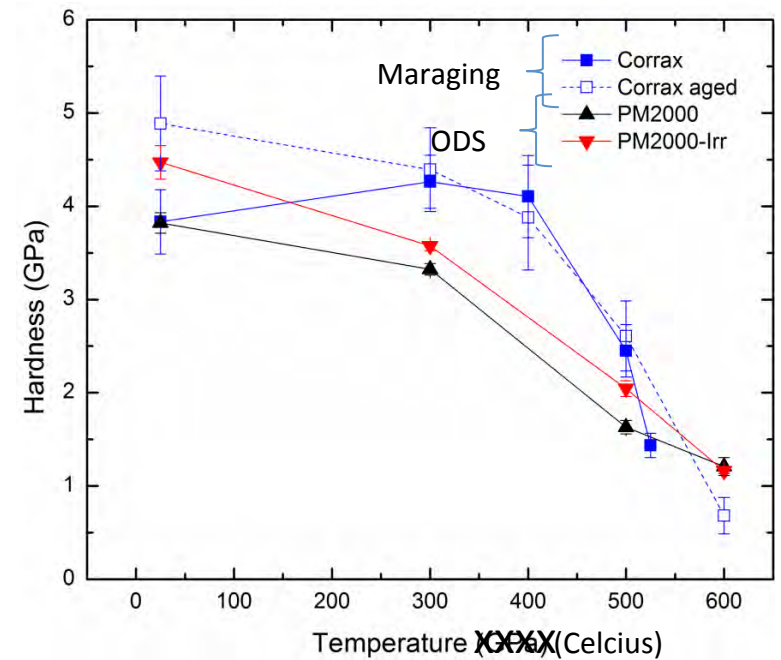
Hardness testing at temperature allows to evaluate the property of the material at operation condition



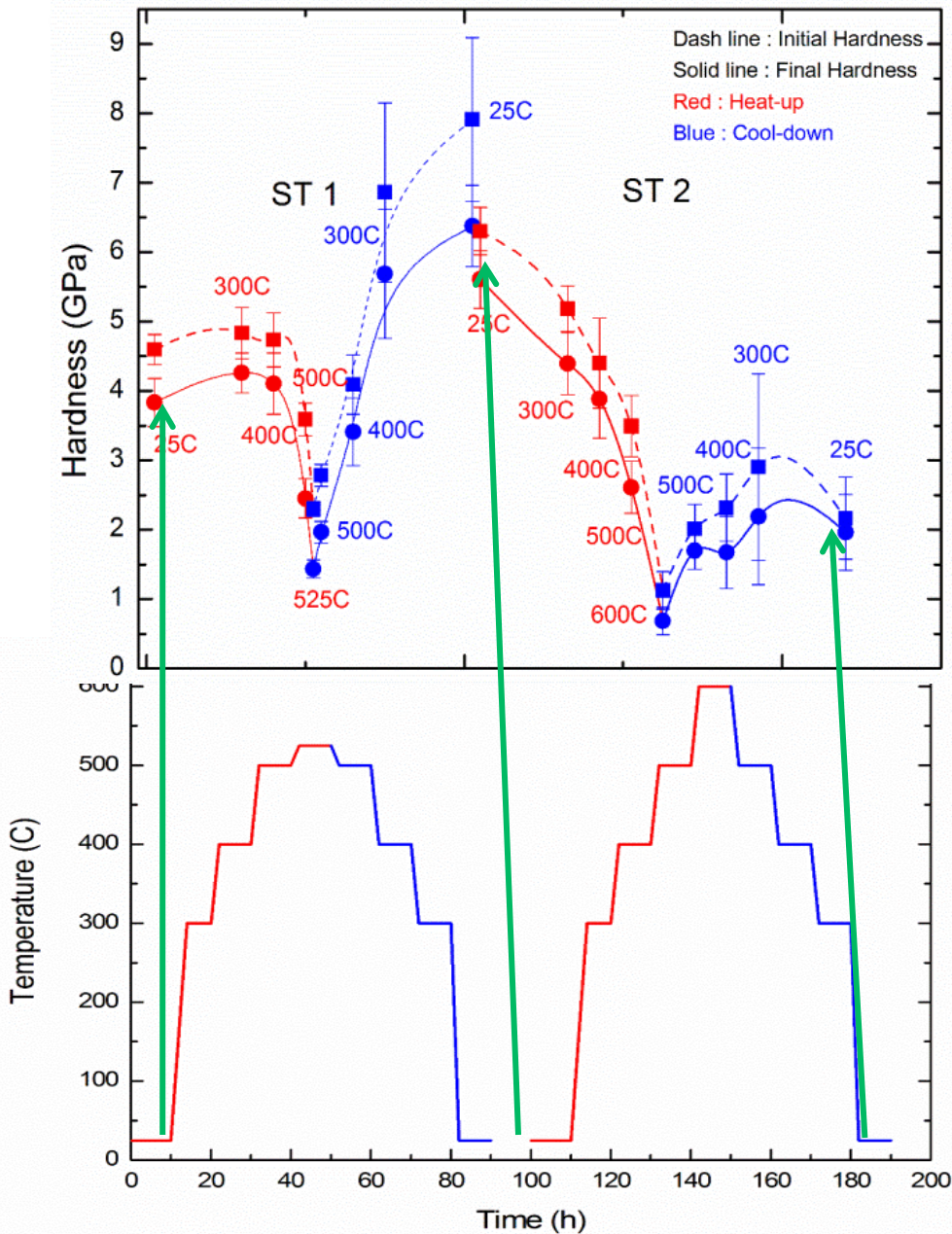
Nanoindeter at UCB



Hardness measurement at temperature



Experimental work performed on heat treatments



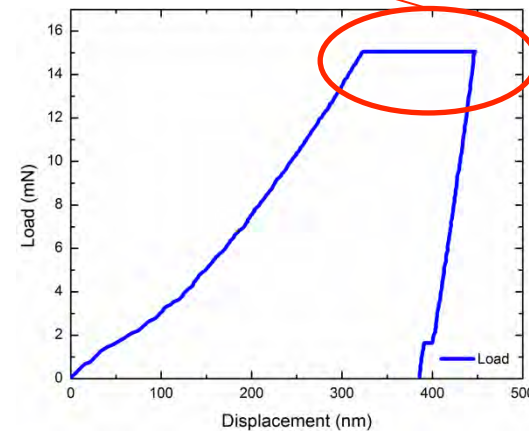
Material softens at temperatures above 400C rapidly.

The hardness at the end of the heat treatment is significantly higher than at the beginning.

The hardness after 600C heat is reduced due to over aging of the material.

Therefore service temperature above 400C is not desirable. Ongoing long term aging study at 400-450C

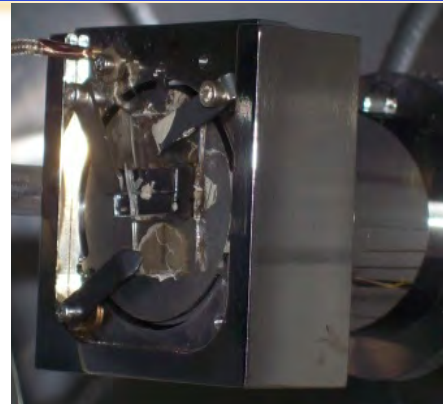
The indentation hardness might also allow the analysis of creep



Experimental work performed on ion beam irradiation



Sample holder paper burn



Sample in beam

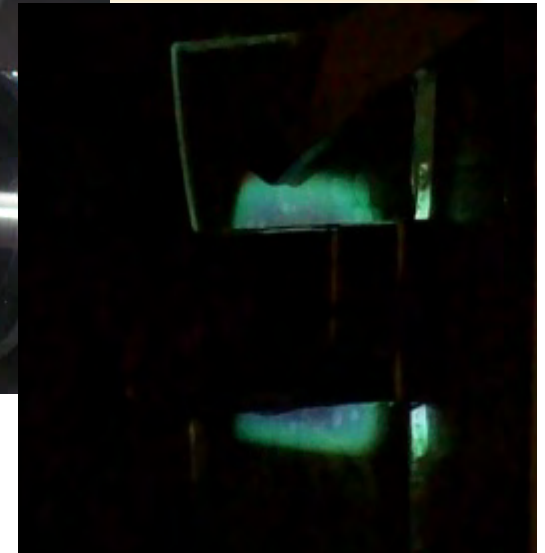
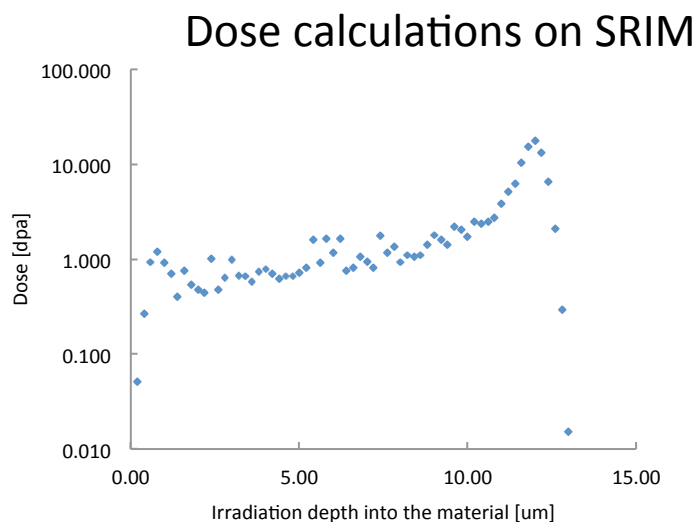


Figure 4: A UCB student refocusing the beam at the IBML facility at LANL

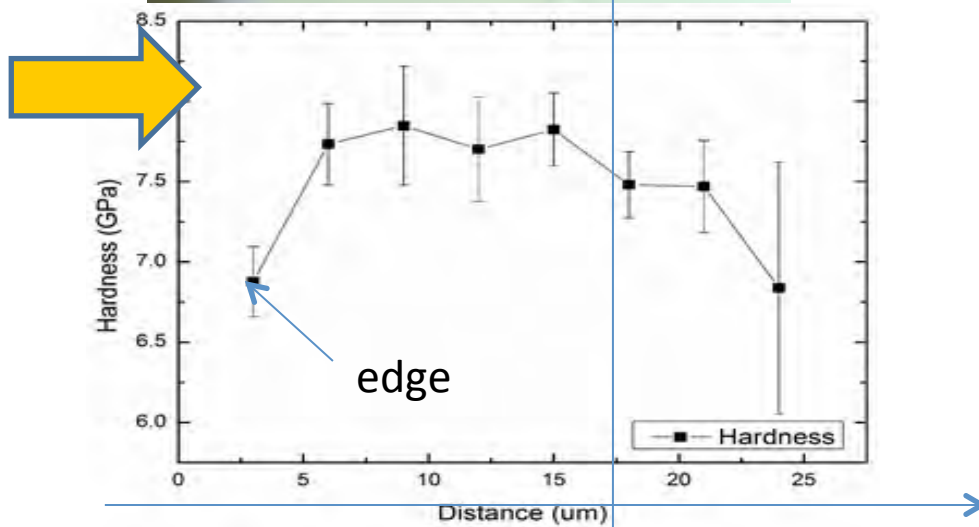
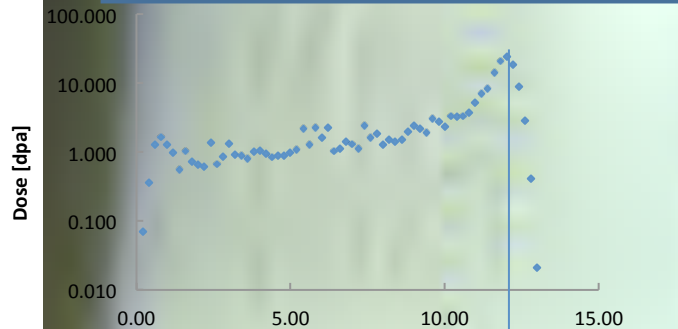
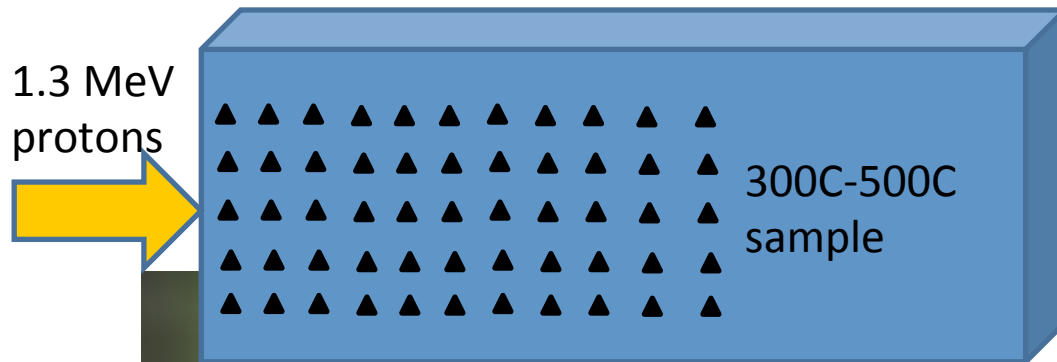


2 irradiations on a total of 4 samples were carried out at this point.

300C 1dpa. However, it was found that the new sample holder designed for high temperature overheated the sample and the irradiation temperature actual was more like 500C.

A second irradiation with 3 different aged materials was conducted to the same dose but with more precise temperature control at 300C.

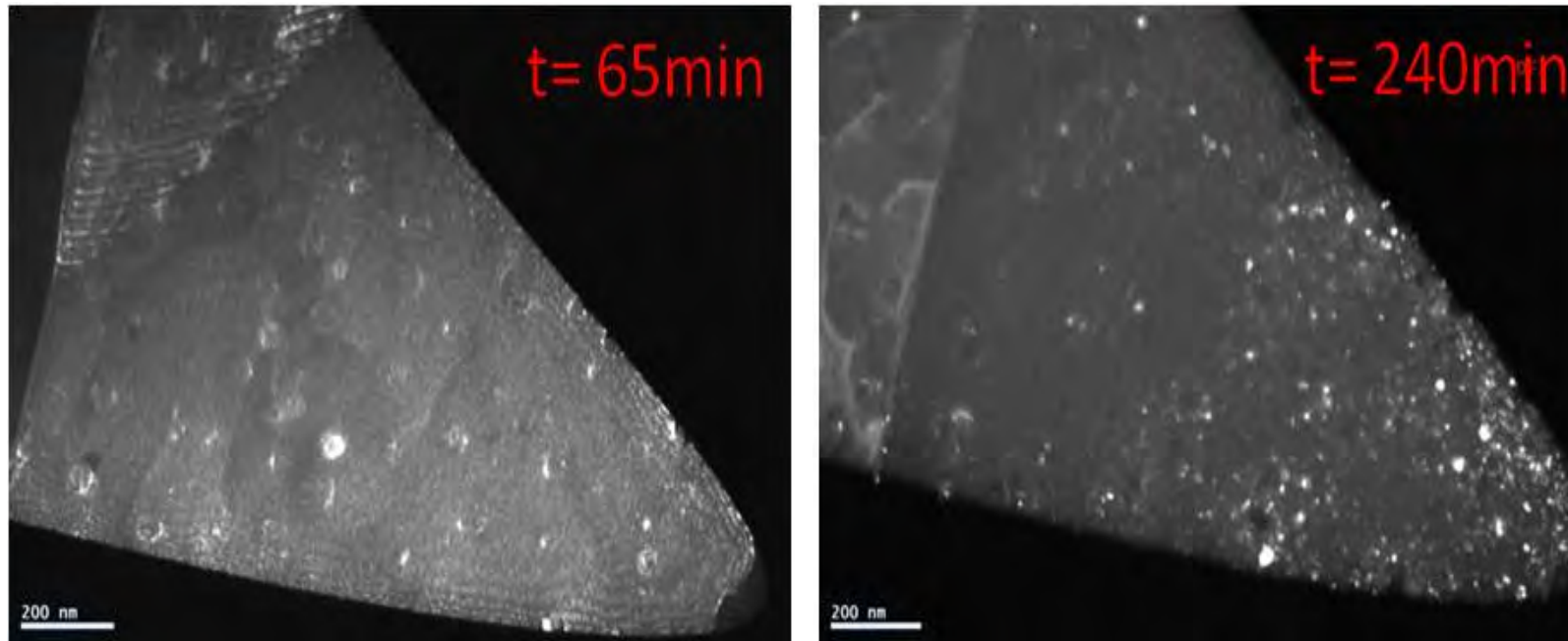
Experimental work performed on ion beam irradiation



It was found that due to the faulty holder the irradiation temperature was not 300C but more likely 500C. The bulk hardness of 7GPa and a subsequent separate temperature calibration experiment confirm this.

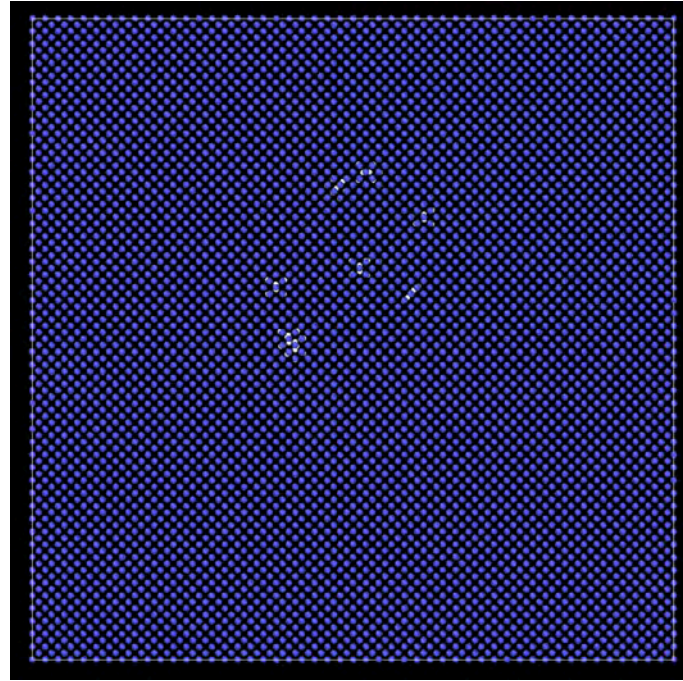
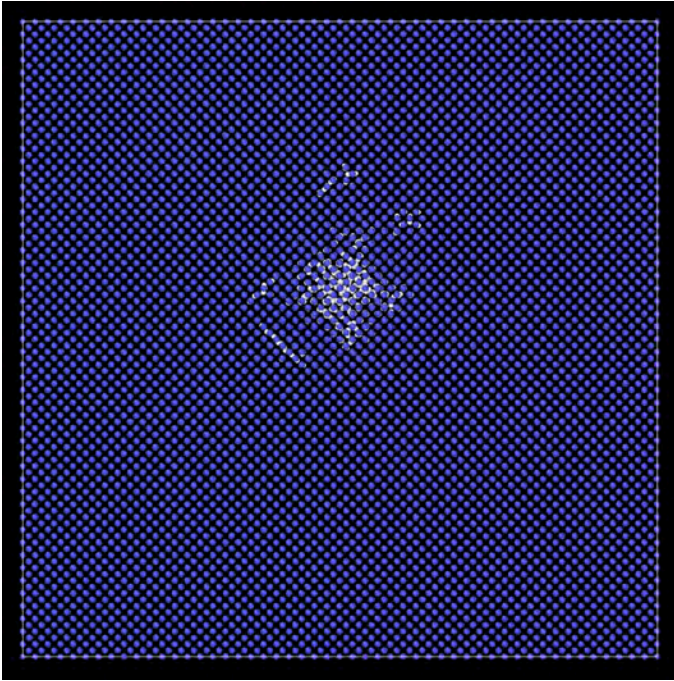
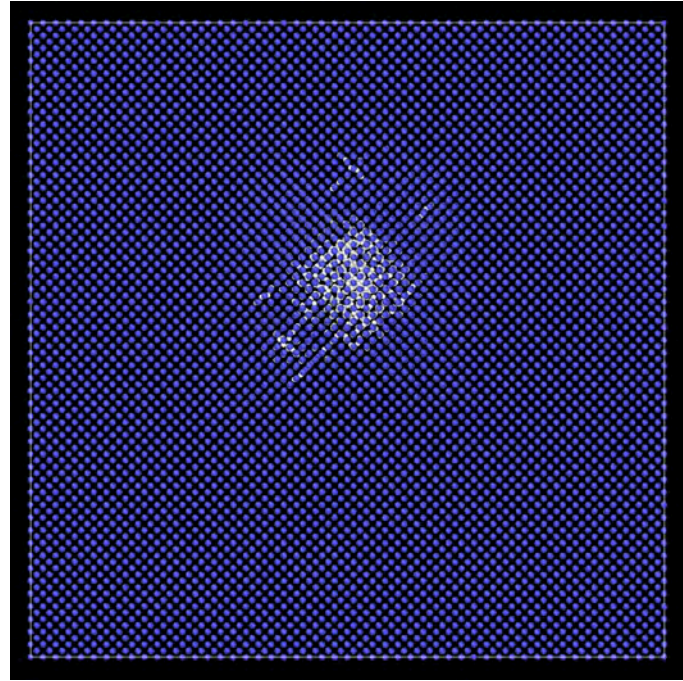
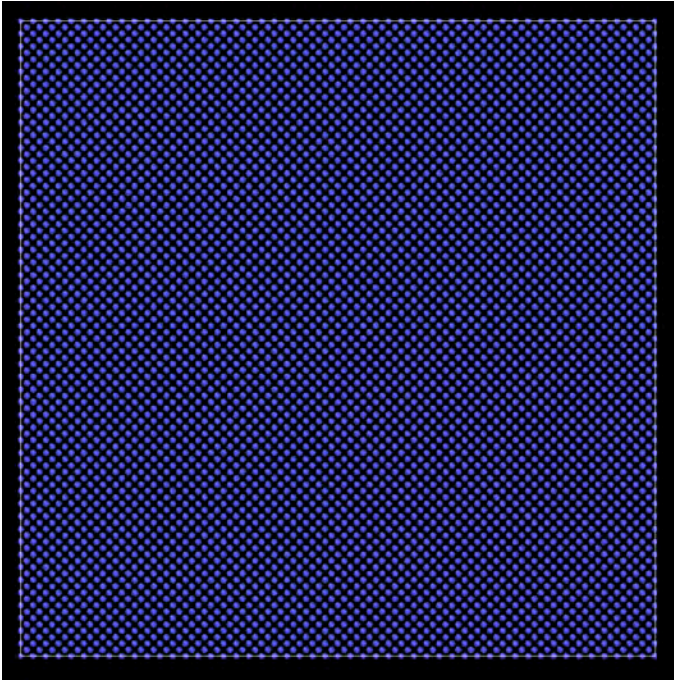
It was found that the 1dpa irradiated zone at ~500C caused a hardening from 3.5GPa to 7.8 GPa. This is close to the peak hardening observed at 525C aging. The bulk hardness stayed at 7GPa. The ion beam effect caused a 0.8 GPa hardness increase.

Experimental work performed on in situ ion beam irradiation at HVEM



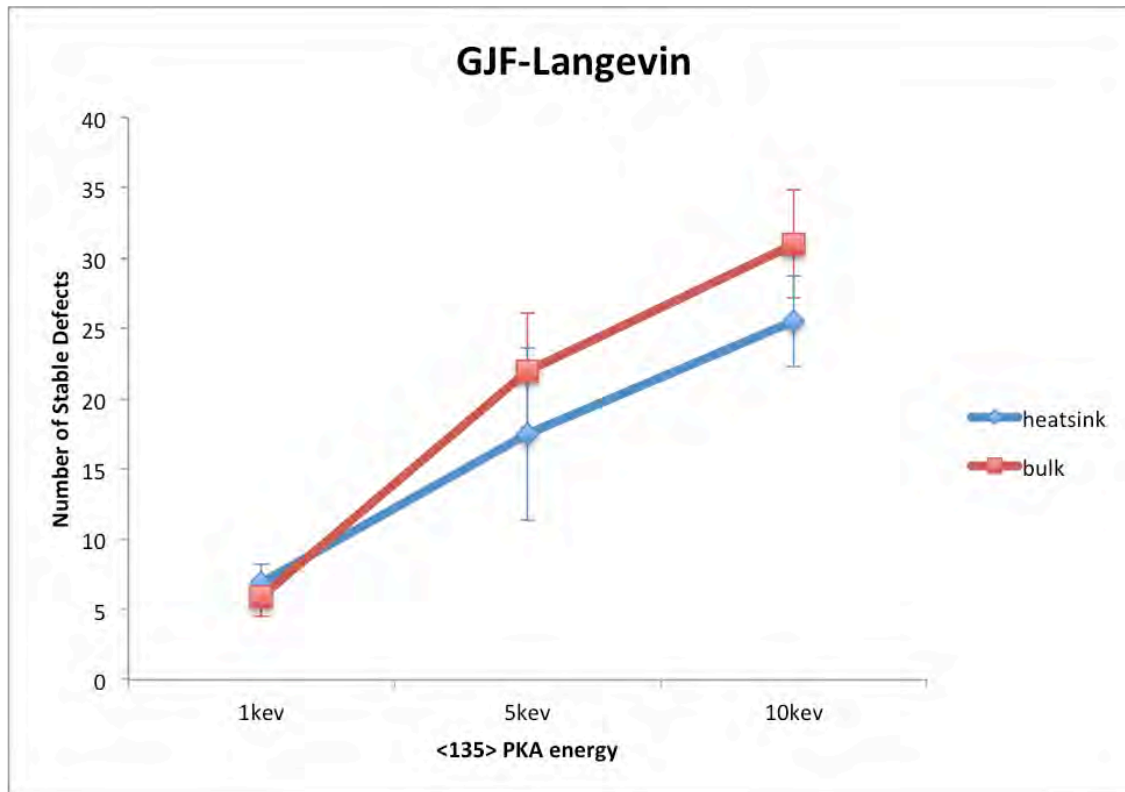
DF-TEM image of the material CORRAX after 65min of irradiation and 240 minutes of irradiation with Kr ions at 500C (exact dose still evaluated). Small white dots (radiation induced precipitates) while other precipitates already pre-exist.

Experiment performed at ANL



Objective of Theory and Simulations

- Electronic structure calculations for the development of Fe-Ni-Al-Fe potentials
- Understand the stability of interfaces and precipitates
Thermal and Radiation studies
- Inclusion of Electronic Stopping and Inelastic Collisions
- ``Byproduct``: New algorithms for general use in Materials Modeling
New, improved Thermostat and Barostat for extended simulations



PKA simulation of long-lived resulting lattice defects with LAMMPS, Including the NEET-developed G-JF Thermostat

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Relevant physical systems – Newtonian dynamics

Self-contained systems:

$$m_i \frac{d^2 r_i}{dt^2} = -\nabla_i U(\{r_j\}) = f_i$$

Ensembles of second order deterministic differential equations:

- Planetary motion, cosmology, Atomic and Molecular Dynamics
- Ballistic motion
- (coupled) oscillator systems, including hyperbolic PDEs

Properties include Energy, (Angular) Momentum conservation.

Noisy systems with hidden degrees of freedom:

$$m_i \frac{d^2 r_i}{dt^2} + \alpha_i \frac{dr_i}{dt} = -\nabla_i U(\{r_j\}) + \eta_i(t) = f_i + \eta_i(t)$$

Ensembles of second order stochastic differential equations:

- Atomic and Molecular systems coupled to, e.g., heat bath
- Diffusive motion
- (coupled) oscillator systems, including hyperbolic PDEs

Properties include equilibrium / Boltzmann distribution.

Discretizing time – Approximating the derivatives

$$v^{n+1} = v^{n+\frac{1}{2}} + \frac{dt}{2}a^{n+\frac{1}{2}} + \frac{dt^2}{8}\dot{a}^{n+\frac{1}{2}} + \mathcal{O}(dt^3)$$

$$v^n = v^{n+\frac{1}{2}} - \frac{dt}{2}a^{n+\frac{1}{2}} + \frac{dt^2}{8}\dot{a}^{n+\frac{1}{2}} + \mathcal{O}(dt^3)$$

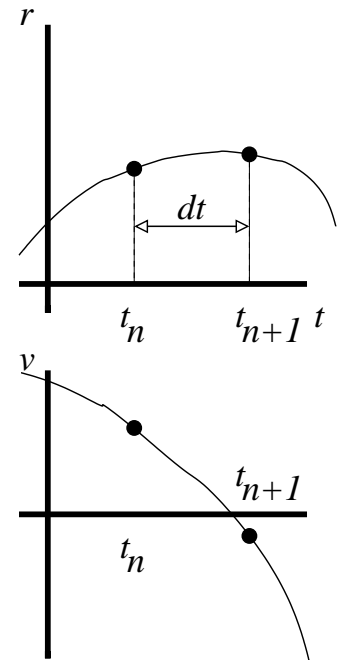
$$\Rightarrow v^{n+1} = v^n + dt a^{n+\frac{1}{2}} + \mathcal{O}(dt^3)$$

$$a^{n+1} = a^{n+\frac{1}{2}} + \frac{dt}{2}\dot{a}^{n+\frac{1}{2}} + \mathcal{O}(dt^2)$$

$$a^n = a^{n+\frac{1}{2}} - \frac{dt}{2}\dot{a}^{n+\frac{1}{2}} + \mathcal{O}(dt^2)$$

$$\Rightarrow a^{n+\frac{1}{2}} = \frac{1}{2}(a^n + a^{n+1}) + \mathcal{O}(dt^2)$$

$$\Rightarrow \underbrace{\begin{cases} r^{n+1} = r^n + dt v^n + \frac{dt^2}{2m} f^n \\ v^{n+1} = v^n + \frac{dt}{2m} (f^n + f^{n+1}) \end{cases}}_{\text{velocity explicit Verlet}} + \mathcal{O}(dt^3)$$



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Noisy systems:

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$$m \frac{d^2 r}{dt^2} + \alpha \frac{dr}{dt} = f(t, r) + \eta(t) \quad , \quad \alpha > 0$$

$$\langle \eta(t) \rangle = 0$$

$$\langle \eta(t) \eta(t') \rangle = 2\alpha k_B T \delta(t - t')$$

Can we fit this into the Verlet method?

What is now important to retain?

$f \equiv 0$:

$$D = \lim_{t \rightarrow \infty} \frac{(r(t) - r(0))^2}{2t} = \frac{k_B T}{\alpha} \quad , \quad \langle T \rangle = \frac{1}{2} m \langle v^2 \rangle = \frac{1}{2} k_B T$$

$$\rho(r) \propto \exp\left(-\frac{U(r)}{k_B T}\right) \Leftrightarrow U_{\text{pmf}}(r) = -k_B T \ln \rho(r)$$

$f = f(r) = -\nabla U(r) \neq 0$:

$$\langle U \rangle = \frac{1}{2} k_B T \quad , \quad \langle T \rangle = \frac{1}{2} m \langle v^2 \rangle = \frac{1}{2} k_B T$$

$$\rho(r) \propto \exp\left(-\frac{U(r)}{k_B T}\right) \Leftrightarrow U_{\text{pmf}}(r) = -k_B T \ln \rho(r)$$

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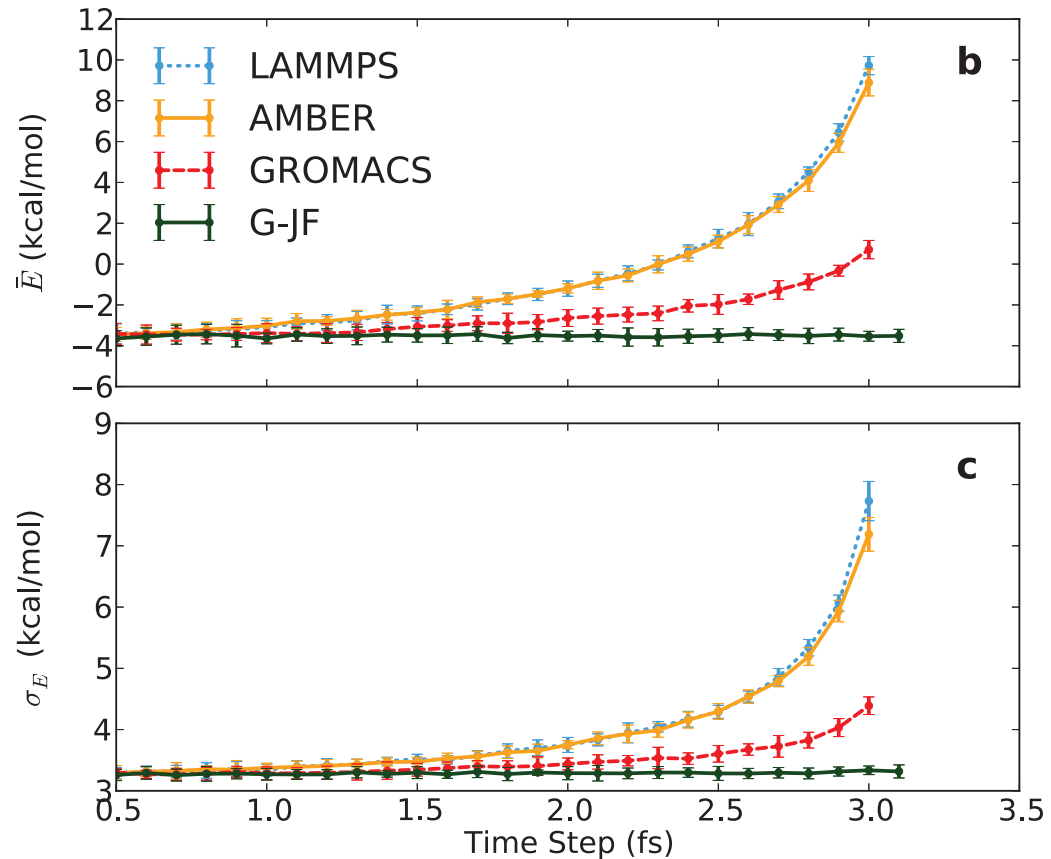
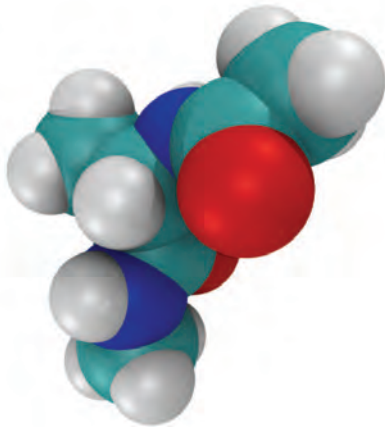
Skeel, Izairre, Mol Phys 100, 3885 (2002);
 Mishra Schlick, J.Chem.Phys 105 299 (1996)
 Melchionna, J.Chem.Phys 127 044108 (2007)

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Applying and Testing the algorithm in existing codes



[Grønbech-Jensen, Hayre, & Farago, (2013) – arXiv:1303.7011]

Now Included into LAMMPS (June 7th, 2013);
 Comparable Barostat in the works;
 Projected inclusion of Electronic Stopping and Inelastic Collisions;
 FeNiAl.....

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Starting Over (Direct Integration):

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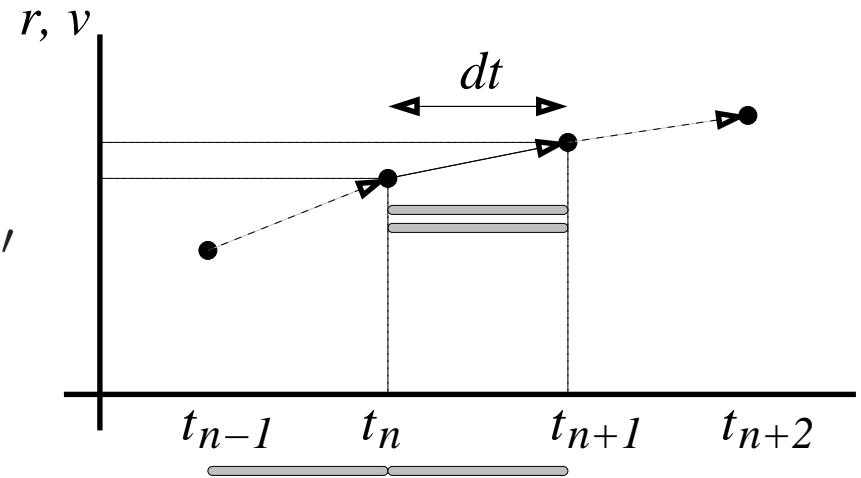
$$m \int_{t_n}^{t_{n+1}} \dot{v} dt' = -\alpha \int_{t_n}^{t_{n+1}} \dot{r} dt' + \int_{t_n}^{t_{n+1}} f dt' + \int_{t_n}^{t_{n+1}} \eta(t') dt'$$

$$v^{n+1} - v^n = \underline{-\frac{\alpha}{m}(r^{n+1} - r^n)} + \frac{1}{m} \int_{t_n}^{t_{n+1}} f dt' + \underline{\frac{dt}{m} \eta^{n+1}}$$

$$\langle \eta^n \rangle = 0$$

$$\langle \eta^n \eta^{n'} \rangle = 2\alpha k_B T \frac{1}{dt} \delta_{n,n'}$$

$$\eta^{n+1} = \frac{1}{dt} \int_{t_n}^{t_{n+1}} \eta(t') dt'$$



$$\int_{t_n}^{t_{n+1}} \dot{r} dt' = \int_{t_n}^{t_{n+1}} v dt'$$

$$r^{n+1} - r^n = \frac{dt}{2}(v^{n+1} + v^n) + \mathcal{O}(dt^3)$$

$$= b dt v^n + \frac{b dt}{2m} \int_{t_n}^{t_{n+1}} f dt' + \frac{b dt^2}{2m} \eta^{n+1}$$

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Thus, ... only the non-dissipative, deterministic force f needs work

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$$r^{n+1} = r^n + b dt v^n + \frac{b dt}{2m} \overbrace{\int_{t_n}^{t_{n+1}} f dt'}^{dt f^n} + \frac{b dt^2}{2m} \eta^{n+1}$$

$$v^{n+1} = v^n - \frac{\alpha}{m} (r^{n+1} - r^n) + \frac{1}{m} \underbrace{\int_{t_n}^{t_{n+1}} f dt'}_{\frac{dt}{2}(f^n + f^{n+1})} + \frac{dt}{m} \eta^{n+1}$$

Which can also be expressed

$$r^{n+1} = r^n + b dt v^n + \frac{b dt^2}{2m} f^n + \frac{b dt^2}{2m} \eta^{n+1}$$

$$v^{n+1} = a v^n + \frac{dt}{2m} (a f^n + f^{n+1}) + \frac{b dt}{m} \eta^{n+1}$$

with

$$a = \frac{1 - \frac{\alpha dt}{2m}}{1 + \frac{\alpha dt}{2m}}, \quad b = \frac{1}{1 + \frac{\alpha dt}{2m}}$$

$$\langle \eta^n \rangle = 0, \quad \langle \eta^n \eta^{n'} \rangle = 2\alpha \frac{k_B T}{dt} \delta_{n,n'}$$

Notice: $\alpha = 0 \Rightarrow$ Störmer-Verlet.

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Result of Analysis[†]

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For all regions within the stability limit, $\Omega_0 dt < 2$:

$$\begin{aligned}\langle U \rangle &= \frac{1}{2} \kappa \langle (r^n)^2 \rangle = \frac{1}{2} k_B T \\ \langle T \rangle &= \frac{1}{2} m \langle (v^n)^2 \rangle = \frac{1}{2} k_B T \left(1 - \frac{\Omega_0^2 dt^2}{4} \right)\end{aligned}$$

i.e., correct spatial Boltzmann distribution regardless of time step, potential curvature, temperature, and friction. Thus, regardless of numerical errors in, e.g., frequency and damping, the system samples the true, continuous time distribution!

Kinetic energy shows the well-known signs of discrete time problems with velocity, re-emphasizing that kinetic energy is a poor measure of temperature in numerical trajectories.

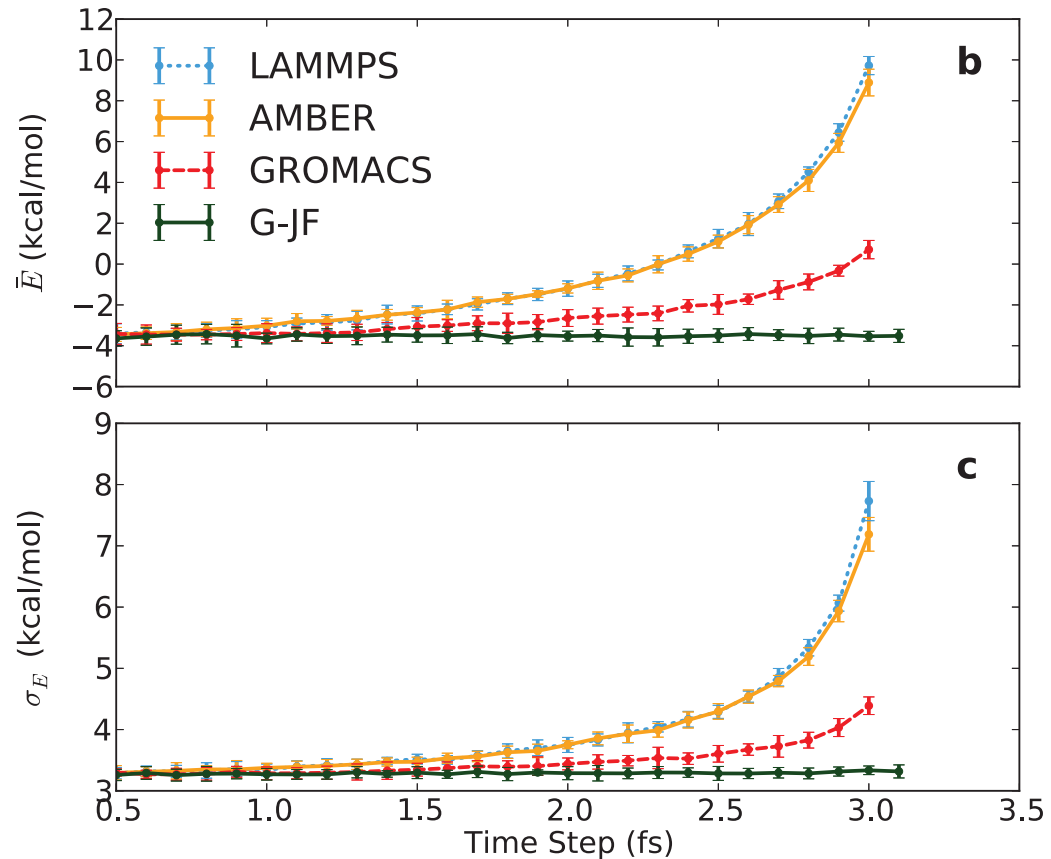
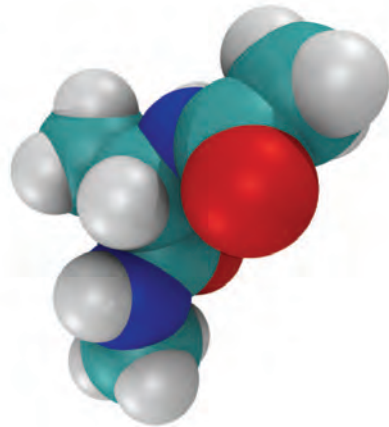
[arXiv:1212.1244/Grønbech-Jensen & Farago, *Molecular Physics* **111**, 983 (2013)]

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- New algorithm implemented at no cost – combining the benefits of Verlet with correct dissipation-fluctuation balance
- Equivalently true for multi-dimensional linear systems; e.g., lattices, PDEs, nonlocal noise. Diagonalization shows that each mode is covered
Now all simulated frequencies respond statistically correct.

- Nonlinear systems.....e.g., comparing PMFs

Time step errors due to deterministic forces remain

Discrete time dissipation-fluctuation balance correct

Tuning α from zero (deterministic) to large is "MD" to "MC"

Large α increasingly accurate even for $\Omega_0 dt \rightarrow 2_-$.

- Do not use kinetic energy measures for temperature; use stochastic method and trust thermodynamics.

[arXiv:1212.1244/Grønbech-Jensen & Farago, *Molecular Physics* **111**, 983 (2013)]

Direction

- Complete ion beam radiation study
 - different temperatures and radiation doses
 - radiation induced coarsening of precipitates
- Complete heat treatment study
 - stability of precipitates
- Helium implantation to trace efficiency of interfaces
- Connect with Modeling (as outlined above)
 - simulations of precipitates (Temp, dose, defects, composition)
 - investigation of new barostat
 - inclusion of bulk friction (ES & Firsov)
- ``Byproduct``: New algorithms for general use in Materials Modeling
 - New, improved Thermostat and Barostat for extended simulations

- Understand operational stability and ``sink-efficiency`` of precipitates